NLO corrections using the Monte Carlo MCFM

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In collaboration with:

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NLO QCD Simulations

- Single top production
 - Fully differential final states see Z. Sullivan's talk Harris, Laenen, Phaf, Sullivan and Weinzierl, 2000
- Diboson production, e.g. $p\bar{p} \to W^+W^- \to \text{leptons}$.
 - Baur et al. lepton correlations only partially included
 Ohnemus, 1994
 Baur, Han and Ohnemus, 1995, 1996
 - Dixon et al. full correlations, anomalous couplings
 Dixon, Kunszt and Signer, 1999
 - MCFM full correlations, singly-resonant contributions

 JC and Ellis, 1999
- Inclusive jets
 - JETRAD 1 and 2 jets only
 Giele, Glover and Kosower, 1993
 - Giele, Kilgore 3 jet production Giele, Kilgore, 2000



NLO QCD Simulations

- Drell-Yan + heavy flavours
 - MCFM $W^{\pm}g^{\star}(\rightarrow b\bar{b})$ Ellis and Veseli, 1998
 - ullet MCFM $Zg^\star(o bar b)$
- Drell-Yan + jets
 - DYRAD handles vector boson + 0 or 1 jets
 Giele, Glover and Kosower, 1993
 - ullet VECBOS handles vector boson + up to 3 (Z)
 - or 4 (W) jets at leading order only

Berends, Kuijf, Tausk and Giele, 1991



MCFM Background

- The Tevatron Run II will be sensitive to processes at the femtobarn level.
- Particularly interesting are final states involving heavy quarks, leptons and missing energy.
- MCFM aims to provide a unified description of such processes at NLO accuracy.
- The extension to NLO is made possible in many cases by the recent calculations of virtual matrix elements involving a vector boson and four partons.
- Similar philosophy, but different approach to Pythia. Whilst Pythia has the advantages of extra radiation (partially included in a NLO calculation) and showering, a fixed order MC may be viewed as theoretically cleaner.



MCFM Process List

Included at NLO

$$egin{aligned} par p &
ightarrow W^\pm/Z & par p &
ightarrow W^+ + W^- \ par p &
ightarrow W^\pm+Z & par p &
ightarrow Z+Z \ par p &
ightarrow W^\pm/Z + H & par p &
ightarrow W^\pm/Z + 1 ext{ jet} \ par p &
ightarrow W^\pm/Z + g^\star \left(
ightarrow bar b
ight) \end{aligned}$$

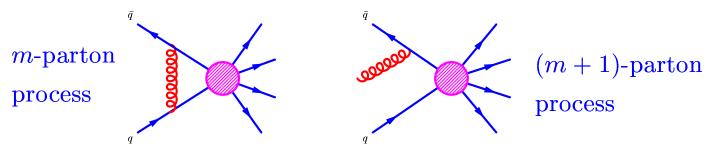
 Various leptonic and/or hadronic decays of the bosons are included as further sub-processes.



No NLO prediction for W/Z+2 jets is available, but this is under construction in MCFM.

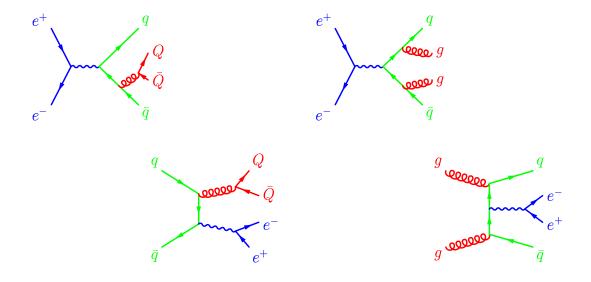
Monte Carlo Ingredients - 1

Helicity amplitudes for the virtual and real ME's



• Many of the NLO matrix elements are obtained by crossing the ones calculated for $e^+e^- \to 4$ jets.

Bern, Dixon, Kosower and Weinzierl, Nucl. Phys. **B489** (1997) 3 Glover and Miller, Phys. Lett. **B396** (1997) 257 Campbell, Glover and Miller, Phys. Lett. **B409** (1997) 503



Monte Carlo Ingredients - 2

- Singular pieces of the real matrix elements must be identified and cancelled by an appropriate set of counter-terms.
- MCFM uses the dipole method to cancel the infrared divergences between real and virtual contributions.

Catani and Seymour, Nucl. Phys. B485 (1997) 291

$$\sigma_{real}^{m+1} = \int_{(m+1)} (d\sigma_{real} - d\sigma_{counter}) + \int_{(m+1)} d\sigma_{counter}$$

$$= (integrable terms) + \sum_{dipoles} \int_{m} d\sigma \otimes \int_{1} dV_{dipole}$$

where the 1-dimensional integral over the dipoles leads to soft and collinear divergences (poles in ϵ).

• These poles manifestly multiply m-parton ME's and may be cancelled against poles from the loop diagrams.

Higgs search using MCFM

ullet Studies using LO Monte Carlos and other event generators show that for a Higgs in the mass range of 100-130 GeV, the most promising channels for discovery at Run II are associated Higgs production.

Stange, Marciano, Willenbrock, Phys. Rev. **D49** (1994) 1354, **D50** (1994) 4491

$$p\bar{p} \longrightarrow W(\to e\nu)H(\to b\bar{b})$$

 $p\bar{p} \longrightarrow Z(\to \nu\bar{\nu}, \ell\bar{\ell})H(\to b\bar{b})$

- Particularly interesting in the light of hints from LEP2.
- Backgrounds for the WH signal:

$$p\bar{p} \longrightarrow W g^{\star}(\to b\bar{b}) \qquad p\bar{p} \longrightarrow t(\to bW^{+})\bar{t}(\to \bar{b}W^{-})$$

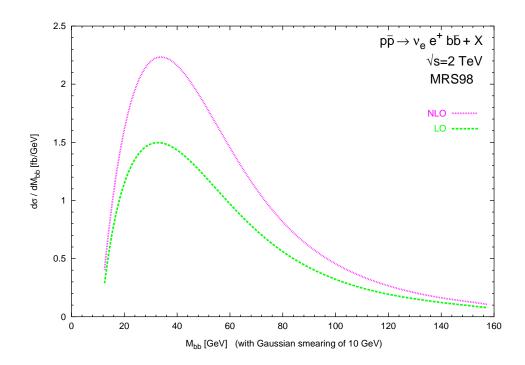
$$p\bar{p} \longrightarrow W Z/\gamma^{\star}(\to b\bar{b}) \qquad p\bar{p} \longrightarrow W^{\pm *}(t(\to bW^{+})\bar{b})$$

$$qg \longrightarrow q't(\to bW^{+})\bar{b}$$



Results for $Wb\bar{b}$

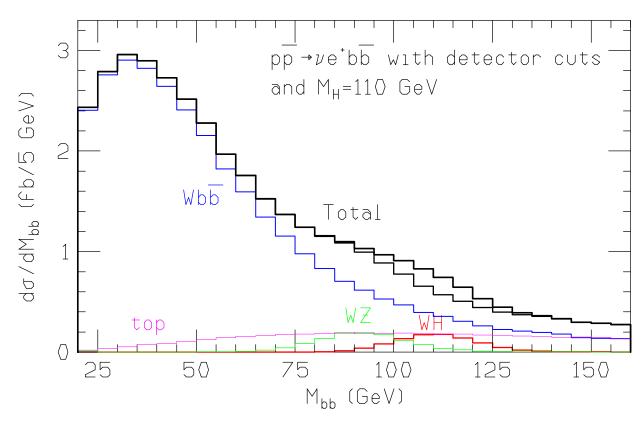
- Use a set of "standard" cuts from the literature, appropriate for the WH study and MRS98 parton distribution functions.
- ullet $m_{bar{b}}$ distribution at LO and NLO, scale of 100 GeV.



ullet The shape changes very little and the K-factor pprox 1.5



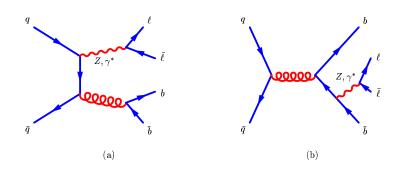
Signal and Backgrounds for $m_H=110~{ m GeV}$

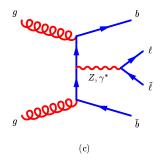


- Double b-tagging efficiency of $\epsilon_{b\bar{b}}=0.45$
- Extraction of the signal requires detailed knowledge of the normalization and the kinematics of the backgrounds.

Results for $Zb\bar{b}$

- ullet New results include radiative corrections, relevant for a further Higgs search in the channel ZH.
- \bullet The required matrix elements are very similar to the $Wb\bar{b}$ case,



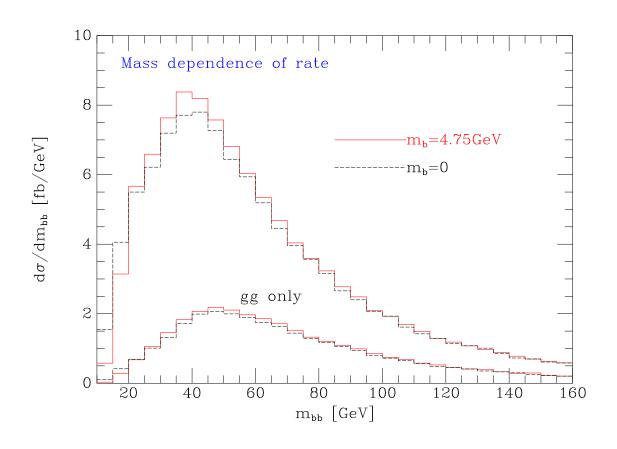


with additional contributions from gg initial states.



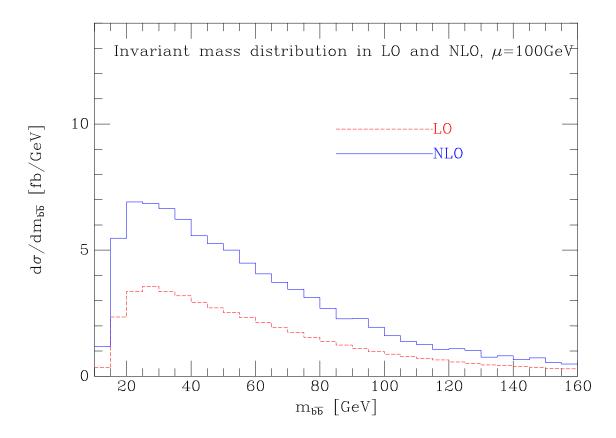
The gg sub-process

• A $b\bar{b}$ pair with a large invariant mass can be produced by the gg initial state process, without off-shell propagators. This gives rise to a large contribution that is important for searches.



$m_{bar{b}}$ mass distribution for $Zbar{b}$

• For a 'conventional' scale of 100 GeV, there is a large K-factor in the region of interest, around 1.8.

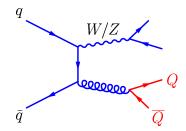


• The entire distribution is changed both in shape and normalization - perhaps suggesting that this scale choice is no longer appropriate (\rightarrow new gg processes).

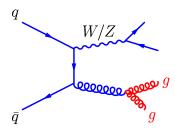


W+2 jets: work in progress

- \bullet View the W+2 jets process as an extension of the $Wb\bar{b}$ and $Zb\bar{b}$ calculations already performed:
 - ullet $Wbar{b}$ part of $qar{q} o W + q'ar{q}'$



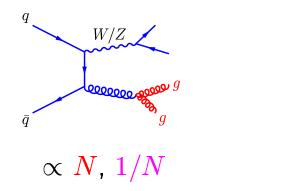
ullet $Zbar{b}$ – contains $gg o Z + qar{q}$ + crossings



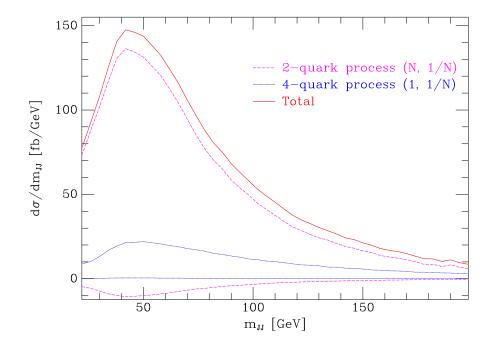
- There are extra parton configurations that we must count.
- The contribution from the diagrams that include real radiation must incorporate the extra singularities due to more configurations of soft/collinear gluons and collinear quark pairs.

W+2 jets at lowest order

• Separate the diagrams by colour structure:



• Typical distribution, standard set of cuts:



W+2 jets: strategy

$$|\mathcal{M}_{NLO}(Vqar{q}gg)|^2 \sim 1 \qquad \longleftarrow ext{Near completion} \ + rac{1}{N^2} \ + rac{1}{N^4}$$

$$|\mathcal{M}_{NLO}(Vqar{q}Qar{Q})|^2 \sim rac{1}{N} \longleftarrow ext{Next target} \ + rac{1}{N^3} \ + rac{1}{N^2} imes \delta_{qQ} \ + rac{1}{N^4} imes \delta_{qQ}$$

ullet Emphasis on W+2 jet first

Conclusions

- Large radiative corrections to the $Wb\bar{b}$ and $Zb\bar{b}$ processes can significantly change estimates of the backgrounds to the processes $p\bar{p}\to WH$ and $p\bar{p}\to ZH$, which will be important search channels at the Tevatron.
- Work is still ongoing in the area of W/Z+2 jet production, for which first results should be available soon.
- MCFM may be downloaded from
 http://www-theory.fnal.gov/people/campbell/mcfm.html